**Table 6.** Summary of the expected neutrino interaction rates in the different detectors for a  $8M_{\odot}$  SN located at 10 kpc (Galactic center). The following notations have been used: IBD, *e*ES and pES stand for Inverse Beta Decay, electron and proton Elastic Scattering, respectively. The final state nuclei are generally unstable and decay either radiatively (notation \*), or by  $\beta^{-}/\beta^{+}$  weak interaction (notation  $\beta^{-,+}$ ). The rates of the different reaction channels are listed, and for LENA they have been obtained by scaling the predicted rates from [65, 66].

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MEMPHYS		LENA		GLACIER	
Interaction	Rates	Interaction	$\operatorname{Rates}$	Interaction	Rates
$ \begin{array}{l} \bar{\nu}_e \ \mathrm{I}\beta \mathrm{D} \\ \nu_e^{(-)CC}(^{16}O, X) \\ \nu_x \ e\mathrm{ES} \end{array} $	$2 \times 10^5$ $10^4$ $10^3$	$ \begin{array}{c} \bar{\nu}_{e} \ \mathrm{I}\beta\mathrm{D} \\ \nu_{x} \ \mathrm{pES} \\ \nu_{x}^{NC}(^{12}C^{*}) \\ \nu_{x} \ e\mathrm{ES} \\ \bar{\nu}_{e}^{CC}(^{12}C, ^{12}B^{\beta^{+}}) \\ \nu_{e}^{CC}(^{12}C, ^{12}N^{\beta^{-}}) \end{array} $	$9 \times 10^{3}$ $7 \times 10^{3}$ $3 \times 10^{3}$ 600 500 85	$\nu_{e}^{CC}({}^{40}Ar, {}^{40}K^{*}) \\ \nu_{x}^{NC}({}^{40}Ar^{*}) \\ \nu_{x} \ e E S \\ \bar{\nu}_{e}^{CC}({}^{40}Ar, {}^{40}Cl^{*}) $	$2.5 \times 10^4$ $3.0 \times 10^4$ $10^3$ 540
Neutronization Burst rates					
LENA	00 70	$\nu_e \text{ eES} / \text{pES}$			
GLACIER	380	$\nu_x^{NC}({}^{40}Ar^*)$			

to all neutrino flavours, would give information on the total flux. GLACIER has also the opportunity to detect  $\nu_e$  by charged-current interactions on  ${}^{40}\text{Ar}$  with a very low energy threshold. The detection complementarity between  $\nu_e$  and  $\bar{\nu}_e$  is of great interest and would assure a unique way of probing the SN explosion mechanism as well as assessing intrinsic neutrino properties. Moreover, the huge statistics would allow spectral studies in time and in energy domain.

We wish to stress that it will be difficult to establish SN neutrino oscillation effects solely on the basis of a  $\bar{\nu}_e$  or  $\nu_e$  spectral hardening, relative to theoretical expectations. Therefore, in the recent literature the importance of model-independent signatures has been emphasized. Here we focus mainly on signatures associated to the prompt  $\nu_e$ neutronization burst, the shock-wave propagation and the Earth matter crossing.

The analysis of the time structure of the SN signal during the first few tens of milliseconds after the core bounce can provide a clean indication if the full  $\nu_e$ burst is present or absent, and therefore allows distinguishing between different mixing scenarios, as indicated by the third column of Tab. 7. For example, if the mass ordering is normal and  $\theta_{13}$  is large, the  $\nu_e$  burst will fully oscillate into  $\nu_x$ . If  $\theta_{13}$  turns out to be relatively large one could be able to distinguish between normal and inverted mass hierarchy.

As discussed above, MEMPHYS is mostly sensitive to the IBD, although the  $\nu_e$  channel can be measured by the elastic scattering reaction  $\nu_x + e^- \rightarrow e^- + \nu_x$  [67]. Of course, the identification of the neutronization burst is the cleanest with a detector exploiting the charged-current absorption of  $\nu_e$  neutrinos, such as GLACIER. Using its unique features of measuring  $\nu_e$  CC events it is possible to probe oscillation physics during the early stage of the SN explosion, while with NC events one can decouple the SN mechanism from the oscillation physics [68, 69].

A few seconds after core bounce, the SN shock wave will pass the density region