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# First events from the CNGS neutrino beam detected in the OPERA experiment 



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#### Abstract

The OPERA neutrino detector at the underground Gran Sasso Laboratory (LNGS) was designed to perform the first detection of neutrino oscillations in appearance mode, through the study of $v_{\mu} \rightarrow v_{\tau}$ oscillations. The apparatus consists of a lead/emulsion-film target complemented by electronic detectors. It is placed in the high-energy, long-baseline CERN to LNGS beam (CNGS) 730 km away from the neutrino source. In August 2006, a first run with CNGS neutrinos was successfully conducted. A first sample of neutrino events was collected, statistically consistent with the integrated beam intensity. After a brief description of the beam and of the various sub-detectors, we report on the achievement of this milestone, presenting the first data and some analysis results.


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## 1. Introduction

The solution of the long-standing solar and atmospheric neutrino puzzles has come from the hypothesis of neutrino oscillations. This implies that neutrinos have nonvanishing and not degenerate masses, and that their flavour eigenstates involved in weak interaction processes are a superposition of their mass eigenstates [1].

Several key experiments conducted in the last few decades with solar neutrinos (see [2] for a review), and with atmospheric, reactor and accelerator neutrinos, have contributed to build-up our present understanding of neutrino mixing. Atmospheric neutrino oscillations, in particular, have been studied by the Super-Kamiokande [3], Kamiokande [4], MACRO [5] and SOUDAN2 [6] experiments. Long baseline experiments confirmed the oscillation hypothesis with accelerator neutrinos: K2K [7] in Japan and MINOS [8] in the USA. The CHOOZ [9] and Palo Verde [10] reactor experiments excluded the $v_{\mu} \rightarrow v_{\mathrm{e}}$ channel as the dominant one in the atmospheric sector.

However, the direct appearance of a different neutrino flavour is still an important open issue. Long-baseline accelerator neutrino beams can be used to probe the atmospheric neutrino signal and confirm the preferred solution of $\nu_{\mu} \rightarrow \nu_{\tau}$ oscillations. In this case, the beam energy should be large enough to produce the heavy $\tau$ lepton. This is one of the main goals of the OPERA experiment [11] that uses the long baseline ( $L=730 \mathrm{~km}$ ) CNGS neutrino beam [12] from CERN to LNGS, the largest underground physics laboratory in the world. The challenge of the experiment is to measure the appearance of $\nu_{\tau}$ from $v_{\mu}$ oscillations. This requires the detection of the short-lived $\tau$ lepton ( $c \tau=87.11 \mu \mathrm{~m}$ ) with high efficiency and low background. The $\tau$ is identified by the detection of its characteristic decay topologies, in one prong (electron, muon or hadron) or in three prongs. The $\tau$ track is measured with a large-mass sampling-calorimeter made of 1 mm thick lead plates (absorber material) inter-spaced with thin emulsion films (highaccuracy tracking devices). This detector is historically called an emulsion cloud chamber (ECC) [11]. Among past applications it was successfully used in the DONUT experiment for the first direct observation of the $\nu_{\tau}$ [13].

The OPERA detector is made of two identical super modules (SMs) each consisting of a target section of about 900 ton made of lead/emulsion-film ECC modules (bricks), of a scintillator tracker detector, needed to pre-localize neutrino interactions within the target, and of a muon spectrometer.

The construction of the CNGS beam has been recently completed and a first run took place in August 2006 with good performance of the facility. First data were collected by the OPERA
detector still without ECC bricks installed, yielding a preliminary measurement of the beam features along with the collection of a number of neutrino interactions (319) consistent with the integrated beam intensity of $7.6 \times 10^{17}$ protons on target (p.o.t.). The OPERA experiment operated very satisfactorily during the run.

## 2. The CNGS beam and the OPERA experiment

The CNGS neutrino beam was designed and optimized for the study of $\nu_{\mu} \rightarrow \nu_{\tau}$ oscillations in appearance mode, by maximizing the number of charged current (CC) $v_{\tau}$ interactions at the LNGS site. A 400 GeV proton beam is extracted from the CERN SPS in $10.5 \mu$ s short pulses with design intensity of $2.4 \times 10^{13}$ p.o.t. per pulse. The proton beam is transported through the transfer line TT41 to the CNGS target T40 [12]. The target consists of a series of thin graphite rods helium-cooled. Secondary pions and kaons of positive charge produced in the target are focused into a parallel beam by a system of two magnetic lenses, called horn and reflector. A 1000 m long decay-pipe allows the pions and kaons to decay into muon-neutrinos and muons. The remaining hadrons (protons, pions, kaons, ...) are absorbed by an iron beam-dump. The muons are monitored by two sets of detectors downstream of the dump; they measure the muon intensity, the beam profile and its centre. Further downstream the muons are absorbed in the rock while neutrinos continue to travel towards Gran Sasso.

The average neutrino energy at the LNGS location is $\sim 17 \mathrm{GeV}$. The $\bar{v}_{\mu}$ contamination is $\sim 4 \%$, the $\nu_{\mathrm{e}}$ and $\bar{\nu}_{\mathrm{e}}$ contaminations are lower than $1 \%$, while the number of prompt $\nu_{\tau}$ from $\mathrm{D}_{\mathrm{s}}$ decay is negligible. The average $\mathrm{L} / \mathrm{E}_{v}$ ratio is $43 \mathrm{~km} \mathrm{GeV}^{-1}$. Due to the Earth's curvature neutrinos from CERN enter the LNGS halls with an angle of about $3^{\circ}$ with respect to the horizontal plane.

Assuming a CNGS beam intensity of $4.5 \times 10^{19}$ p.o.t. per year and a five year run about 31000 CC plus neutral current (NC) neutrino events will be collected by OPERA from interactions in the lead-emulsion target. Out of them 95 (214) CC $v_{\tau}$ interactions are expected for oscillation parameter values $\Delta m_{23}^{2}=2 \times 10^{-3} \mathrm{eV}^{2}\left(3 \times 10^{-3} \mathrm{eV}^{2}\right)$ and $\sin ^{2} 2 \theta_{23}=1$. Taking into account the overall $\tau$ detection efficiency the experiment should gather $10-15$ signal events with a background of less than one event.

In the following, we give a brief description of the main components of the OPERA detector.
Each of the two SMs (SM1 and SM2) consists of 103168 lead/emulsion bricks arranged in 31 target planes (figure 1), each one followed by two scintillator planes with an effective granularity of $2.6 \times 2.6 \mathrm{~cm}^{2}$. These planes serve as trigger devices and allow selecting the brick containing a neutrino interaction. A muon spectrometer at the downstream end of each SM allows to measure the muon charge and momentum. A large size anti-coincidence detector placed in front of SM1 allows to veto (or tag) interactions occurring in the material and in the rock upstream of the target.

The construction of the experiment started in spring 2003. The first instrumented magnet was completed in May 2004 together with the first half of the target support structure. The second magnet was completed in the beginning of 2005. In spring 2006, all scintillator planes were installed. The production of the ECC bricks started in October 2006 with the aim of completing the full target for the high-intensity run of 2007.

The run of August 2006 was conducted with electronic detectors only, taking neutrino interactions in the rock upstream of the detector, in the passive material of the mechanical


Figure 1. Schematic drawing of the OPERA detector at LNGS.
structure and in the iron of the spectrometers. In addition, the information from a tracking plane made of pairs of emulsion films (changeable sheets, CS) was used to study the association between emulsion-film segments with tracks reconstructed in the target tracker (TT). Figure 2 shows a photograph of the detector in the underground Hall C of LNGS as it was during the neutrino run.

### 2.1. The electronic detectors

The needs of adequate spatial resolution for high brick finding efficiency, for good calorimetric measurement of the events, as well as the requirement of covering large surfaces ( $\sim 6000 \mathrm{~m}^{2}$ ), impose strong requirements on the TT. Therefore, the cost-effective technology of scintillating strips with wavelength shifting fibre readout was adopted.

The polystyrene scintillator strips are 6.86 m long, 10.6 mm thick and 26.3 mm wide. A groove in the centre of the strip houses the 1 mm diameter fibre. Multi-anode, 64-pixel photomultipliers are placed at both ends of the fibres. A basic unit of the TT called module consists of 64 strips glued together. One plane of four modules of horizontal strips and one of four modules of vertical strips form a scintillator wall providing $X-Y$ track information. The readout electronics is based on a 32 -channel ASIC [14] that outputs a charge proportional to the signal delivered by each pixel of the photomultipliers with a dynamic range from 1 to 100 photoelectrons.

Muon identification and charge measurement are needed for the study of the muonic $\tau$-decay channel and for the suppression of the background from the decay of charmed particles, featuring the same topology. Each muon spectrometer [15] consists of a dipolar magnet made of two iron arms for a total weight of 990 ton. The measured magnetic field intensity is 1.52 T . The two arms are interleaved with vertical, 8 m long drift-tube planes for the precise measurement


Figure 2. Photograph of OPERA in the LNGS Hall C.
of the muon-track bending. Planes of resistive plates chambers (RPCs) are inserted between the iron plates of the arms, providing a coarse tracking inside the magnet, range measurement of the stopping particles and a calorimetric analysis of hadrons.

In order to measure the muon momenta and determine their sign with high accuracy, the precision tracker (PT) is built of thin walled aluminum tubes with 38 mm outer diameter and 8 m length [16]. Each of the $\sim 10000$ tubes has a central sense wire of $45 \mu \mathrm{~m}$ diameter. They can provide a spatial resolution better than $300 \mu \mathrm{~m}$. Each spectrometer is equipped with six 4 -fold layers of tubes.

RPCs identify penetrating muons and measure their charge and momentum in an independent way with respect to the PT. They consist of electrode plates made of 2 mm thick plastic laminate of high resistivity painted with graphite. Induced pulses are collected on two pickup strip planes made of copper strips glued on plastic foils placed on each side of the detector. The number of individual RPCs is 924 for a total detector area of $3080 \mathrm{~m}^{2}$. The total number of digital channels is about 25000 , one for each of the 2.6 cm (vertical) and 3.5 cm (horizontal) wide strips.

In order to solve ambiguities in the track spatial-reconstruction each of the two drift-tube planes of the PT upstream of the dipole magnet is complemented by an RPC plane with two $42.6^{\circ}$ crossed strip-layers called XPCs. RPCs and XPCs give a precise timing signal to the PTs.

Finally, a detector made of glass RPCs is placed in front of the first SM, acting as a veto system for interactions occurring in the upstream rock. The veto detector was not yet operational for the August 2006 run. The PT was in the commissioning phase with two working planes. The TT and the RPCs already passed a full commissioning with cosmic-ray muons before the run ${ }^{38}$.

OPERA has a low data rate from events due to neutrino interactions well localized in time, in correlation with the CNGS beam spill. The synchronization with the spill is done offline via GPS. The detector remains sensitive during the inter-spill time and runs in a trigger-less mode. Events detected out of the beam spill (cosmic-ray muons, background from environmental radioactivity, dark counts) are used for monitoring. The global DAQ is built as a standard Ethernet network whose 1147 nodes are the Ethernet controller mezzanines plugged on controller boards interfaced to each sub-detector specific front-end electronics. A general 10 ns clock synchronized with the local GPS is distributed to all mezzanines in order to insert a time stamp to each data block. The event building is performed by sorting individual subdetector data by their time stamps.

### 2.2. Emulsion films, bricks and related facilities

An R\&D collaboration between the Fuji Company and the Nagoya group allowed the large scale production of the emulsion films needed for the experiment (more than 12 million individual films) fulfilling the requirements of uniformity of response and of production, time stability, sensitivity, schedule and cost [17]. The main peculiarity of the emulsion films used in high energy physics compared to normal photographic films is the relatively large thickness of the sensitive layers ( $\sim 44 \mu \mathrm{~m}$ ) placed on both sides of a $205 \mu \mathrm{~m}$ thick plastic base.

A target brick consists of 56 lead plates of 1 mm thickness and 57 emulsion films. The plate material is a lead alloy with a small calcium content to improve its mechanical properties. The transverse dimensions of a brick are $12.7 \times 10.2 \mathrm{~cm}^{2}$ and the thickness along the beam direction is 7.5 cm (about 10 radiation lengths). The bricks are housed in support structures placed between consecutive TT walls.

In order to reduce the emulsion scanning load the use of changeable sheets, successfully applied in the CERN CHORUS experiment [18], was extended to OPERA. CS doublets are attached to the downstream face of each brick and can be removed without opening the brick. Charged particles from a neutrino interaction in the brick cross the CS and produce a trigger in the TT scintillators. Following this trigger the brick is extracted and the CS developed and analysed in the scanning facility at LNGS. The information of the CS is used for a precise prediction of the position of the tracks in the most downstream films of the brick, hence guiding the so-called scan-back vertex-finding procedure.

The hit brick finding is one of the most critical operations for the success of the experiment, since one aims at high efficiency and purity in detecting the brick containing the neutrino interaction vertex. This requires the combination of adequate precision of the TT, precise extrapolation and high track finding efficiency in the CS scanning procedure. During the neutrino run of August 2006 a successful test of the whole procedure was performed by using an emulsion

[^1]detector plane consisting of a matrix of $15 \times 20$ individual CS doublets with overall transverse dimensions of $158 \times 256 \mathrm{~cm}^{2}$ inserted in one of the SM2 target planes.

The construction of more than 200000 bricks for the neutrino target is accomplished by an automatic machine, the brick assembly machine, operating underground in order to minimize the number of background tracks from cosmic-rays and environmental radiation. Two brick manipulating systems on the lateral sides of the detector position the bricks in the target walls and also extract those bricks containing neutrino interactions.

While running the experiment, after the analysis of their CS doublets, bricks with neutrino events are brought to the LNGS external laboratory, exposed for several hours to cosmic-ray muons for film alignment [19] and then disassembled. The films are developed with an automatic system in parallel processing chains and dispatched to the scanning labs.

The expected number of bricks extracted per running-day with the full target installed and CNGS nominal intensity is about 30 . The large emulsion surface to be scanned requires fast automatic microscopes continuously running at a speed of $\sim 20 \mathrm{~cm}^{2}$ film surface per hour. This requirement has been met after R\&D studies conducted using two different approaches by some of the European groups of the Collaboration (ESS) [20] and by the Japanese groups (S-UTS) [21].

## 3. The first run with CNGS neutrinos

For a detailed description of the CNGS beam operation during the first run with neutrinos of August 2006 we refer to the official CNGS WEB page [12]. The commissioning of the beam started on 10 July 2006 following a series of technical tests of individual components performed from February to May. During this phase the SPS delivered $7 \times 10^{15}$ p.o.t., equivalent to 1 h of CNGS running with nominal intensity.

The first shot of the extracted proton beam onto the CNGS target was made on 11 July. A low intensity run with neutrinos took then place from 18 to 30 August 2006 with a total integrated intensity of $7.6 \times 10^{17}$ p.o.t. (figure 3 ). The beam had been active for a time equivalent to about five days. The low intensity was partly due to the chosen SPS cycle and to the intensity of the spill that was $55 \%$ of the nominal value during the first part of the run and $70 \%$ during the second part.

The GPS clock used to synchronize the CERN accelerators and OPERA had been finetuned before the start of data-taking. At CERN the current pulse of the kicker magnet used for the beam extraction from the SPS to the TT41 line was time-tagged by a GPS unit with absolute time (UTC) calibration. An analogous GPS at the LNGS site provided the UTC timing signal to OPERA. The resulting accuracy in the time synchronization between CERN and OPERA timing systems was better than 100 ns . However, during the first days of the run a time offset of $100 \mu \mathrm{~s}$ was observed due to problems in adjusting the time tagging of the kicker pulse. This offset was eventually reduced to 600 ns .

The OPERA detector started collecting neutrino interactions from the very first beam spills with nearly all electronic detectors successfully operating. Altogether, 319 neutrino events, with an estimated $5 \%$ systematic uncertainty, were taken by OPERA during the August run. This is consistent with the 300 expected events for the given integrated intensity of $7.6 \times 10^{17}$ p.o.t.. The analysis of the CNGS data conducted at CERN and the comparison with simulations is in progress. Once completed, we expect to reach a $20 \%$ systematic error on the prediction


Figure 3. The CNGS proton integrated intensity during the August 2006 run.
of the number of muon events from neutrino interactions in the rock. This error is due to uncertainties in the neutrino flux prediction, in the cross-section and in the muon transport in the rock.

The event analysis was performed in two ways. In the first one the event timing information was treated as a basic selection tool, since the time window of beam events is well sized in a $10.5 \mu$ s interval, while the uniform cosmic-ray background corresponded to $10^{-4}$ of the collected statistics (figure 4). The second analysis dealt with the reconstruction of track-like events disregarding timing information. Neutrino events are classified as: (i) CC neutrino interactions in the rock upstream of the detector or in the material present in the hall leading to a penetrating muon track (figure 5, top-left); (ii) CC and NC neutrino interactions in the target material (figure 5, top-right and bottom-right) and CC interactions in the iron of the spectrometers (figure 5, bottomleft).

The $\theta$ angular distribution with respect to the horizontal axis obtained by selecting singletrack events is shown in figure 6. Events were selected with a minimum number of six layers of fired RPCs in each spectrometer. In the same figure, the distribution of simulated cosmic-ray muons from [5] is also shown. The comparison between experimental data and Monte Carlo events proved the beam-induced nature of the muons in the peak around the horizontal direction. By counting events selected with topological criteria we found $\sim 10 \%$ of the events corresponding to beam spill data missing in the CERN database. A Gaussian fit to the $\theta$ angle of the events on-time with the beam (shown in the inset of figure 6) yielded a mean muon angle of $3.4 \pm 0.3^{\circ}$ in agreement with the value of $3.3^{\circ}$ expected for neutrinos originating from CERN and travelling under the earth surface to the LNGS underground halls. The systematic error on $\theta$ was evaluated to be $3 \%$ by applying different track reconstruction models.


Figure 4. Time distribution of events collected in the neutrino run. The event time difference with respect to the closest extraction is shown in the right histogram.

Concerning the analysis of the CS detector, 9 muons produced by neutrino interactions in the rock surrounding the detector crossed the CS plane surface. Five muon tracks predicted by the electronic detectors were found by scanning the emulsion films. The reasons for inefficiency can be traced back to the tight cuts applied in this preliminary analysis and in the significant decrease of the fiducial volume. In fact, the dead space between adjacent emulsion films was $\sim 10 \%$ and the scanning was only performed up to 3 mm from the film edge, bringing the overall dead space to $\sim 20 \%$. However, the test proved the capability in passing from the centimetre scale of the electronic tracker resolution to the micrometric resolution of nuclear emulsions. The angular difference between predicted and found tracks is better than 10 mrad , largely dominated by the electronic detector resolution. Figure 7 shows the display of one of the 5 reconstructed events.

## 4. Conclusions

We reported the first detection of neutrino events from the long baseline CERN CNGS beam with the OPERA experiment in the underground Gran Sasso Laboratory. The electronic detectors of the experiment performed successfully with an overall data-taking efficiency larger than $95 \%$ during the August 2006 run. The scintillator TT and the spectrometers equipped with RPCs allowed to identify muon tracks from CC neutrino interactions occurring in the rock and in the material upstream of the detector, as well as in the detectors themselves.


Figure 5. Display of neutrino events from the CNGS run. For each event the top and side views are shown, respectively. The SM targets are indicated in blue, the spectrometers in light brown, TT and RPC hits in red. See the text for event classification.

Three hundred and nineteen neutrino-induced events were collected for an integrated intensity of $7.6 \times 10^{17}$ p.o.t. in agreement with the expectation of 300 events. The reconstructed zenith-angle distribution from penetrating muon tracks is centred at $3.4^{\circ}$ with a $10 \%$ statistical error, as expected for neutrinos originating from CERN and travelling under the earth surface to LNGS.

A test of the association between muon tracks reconstructed with the electronic detectors and with an emulsion detector plane was also successfully performed, proving the capability of passing from the centimetre scale of the electronic tracker resolution to the micrometric resolution of nuclear emulsions. The angular difference in the track association is better than 10 mrad , largely dominated by the electronic detector resolution.

The success of this first OPERA run with CNGS neutrinos is the first step towards the operation of the complete detector.


Figure 6. Angular distribution of beam-induced and cosmic-muon events taken with the electronic detectors (black points). The histogram indicates the predictions from cosmic-ray simulations [5]. The inset shows the angular distribution of on-time beam events.


Figure 7. Left: display of one event with the muon passing through the CS detector plane. Only hits of the electronic detectors close to the CS plane are shown; the vertical segment indicates the position of the CS doublet intercepted by the track. Right: display of the corresponding four micro-tracks reconstructed in the CS doublet.

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[^1]:    ${ }^{38}$ The cosmic muon flux in the LNGS Hall C integrated over the full solid angle is about 1 muon $\mathrm{m}^{-2} \mathrm{~h}^{-1}$.

