

# Physics Potential of the SPL Super Beam

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**Abstract.** Performances of a neutrino beam generated by the CERN SPL proton driver are computed considering a 440 kton water Čerenkov detector at 130 km from the target.  $\theta_{13}$  sensitivity down to  $1.2^\circ$  and a  $\delta$  sensitivity comparable to a Neutrino Factory, for  $\theta_{13} \geq 3^\circ$ , are within the reach of such a project.

## 1. Introduction

The planned Super Proton Linac (SPL) is a 2.2 GeV proton beam of 4 MW power [1] working with a repetition rate of 75 Hz delivering  $1.5 \cdot 10^{14}$  protons per pulse ( $10^{23}$  protons on target (pot) in a  $10^7$  s conventional year). It could be the first stage of a CERN based Neutrino Factory or of a Beta Beam.

Studies of the capabilities of a neutrino beam generated by SPL have been already published in [2], [3], in this paper the fluxes and the overall physical performances will be reviewed in the light of the new design of the beam optics, specially optimized for the SuperBeam needs [4]. They are computed for a gigantic water Čerenkov detector, as the proposed UNO detector [5] (440 kton fiducial) located in the Modane laboratory under the Frejus tunnel at a baseline of 130 km from CERN.

The SPL SuperBeam capabilities would constitute a natural follow-up of the JHF phase I experiment [6], with excellent sensitivity on  $\theta_{13}$  (section 3) and good sensitivity on the CP phase  $\delta$  (section 4). Furthermore the SPL SuperBeam could be used to complement the results of a Neutrino Factory experiment, helping in resolving the ambiguities, as discussed in [7]; or could be combined with a Beta Beam, as discussed in [8].

Signal efficiency and backgrounds have already been discussed in [2], they are computed by using the NUANCE neutrino event generator [9] and reconstructing the events with standard SuperKamiokande algorithms, with the addition of improved  $\pi^0$  rejection algorithms. They can be summarized as signal efficiency  $\epsilon \simeq 70\%$  and  $\pi^0$  and  $\mu/e$  background rejection, normalized to the non oscillated  $\nu_\mu$  charged current interactions,  $f_B^{\pi^0} = 4.2 \cdot 10^{-4}$ ,  $f_B^{\mu/e} = 3 \cdot 10^{-3}$ .

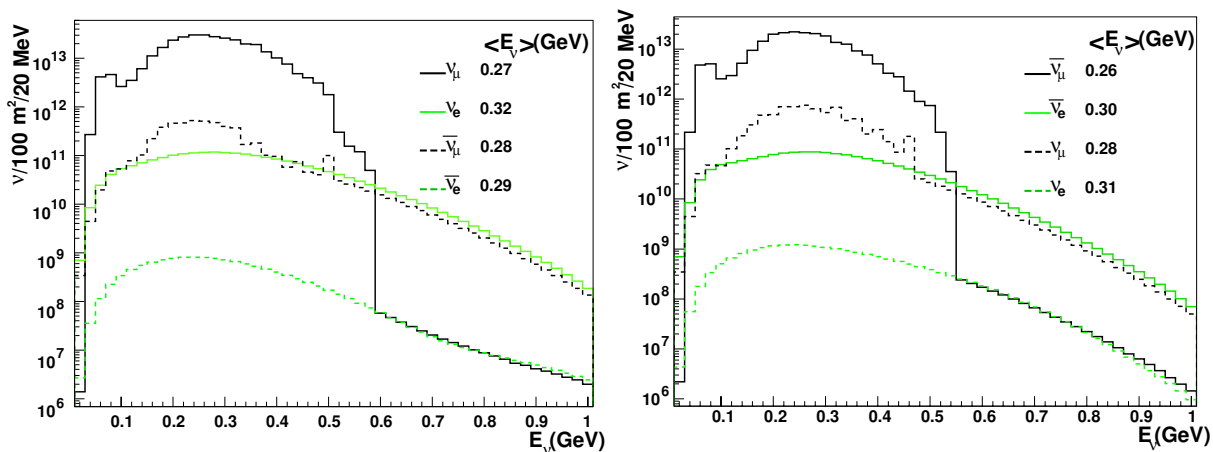
## 2. Fluxes

Details of the new beam optics can be found in [4]. The use of an horn and a reflector increases by  $\sim 40\%$  the overall  $\nu_\mu$  flux with respect to the former single horn optics, slightly increases the  $\nu_e$  contamination, while the  $\bar{\nu}_\mu$  and  $\bar{\nu}_e$  contaminations are reduced by  $\sim 30\%$ .

The length of the decay tunnel has been re-optimized having in mind CP searches more than  $\theta_{13}$ . Table 1 reports details of the beam properties as function of the length of the decay tunnel, including the sensitivity on  $\theta_{13}$  for a 2200 kton/year exposure. In spite of the fact that the  $\theta_{13}$  sensitivity is maximum for the lowest length (20 m), a 60 m decay length is preferred because of the lower  $\bar{\nu}_\mu$  contamination, that results in a better CP sensitivity. The neutrino spectra for the  $\pi^+$  and  $\pi^-$  focussed beams are displayed in Fig. 1.

**Table 1.** Neutrino fluxes and contamination for different values of the decay tunnel length. The last line refers of the single horn optics of ref. [2].

Length (m)	$\pi^+$ focus			$\pi^-$ focus			$\pi^+$ focus
	$\nu_\mu$ ( $\nu/m^2/yr$ ) (@50 km)	$\nu_e$ (%)	$\bar{\nu}_\mu$ (%)	$\bar{\nu}_\mu$ ( $\nu/m^2/yr$ ) (@50 km)	$\bar{\nu}_e$ (%)	$\nu_\mu$ (%)	$\theta_{13}$ (90%CL) (2200 kton/yr)
20	$2.43 \cdot 10^{+12}$	0.38	1.71	$1.73 \cdot 10^{+12}$	0.41	3.9	1.20
60	$3.23 \cdot 10^{+12}$	0.67	1.50	$2.25 \cdot 10^{+12}$	0.70	3.3	1.25
100	$3.35 \cdot 10^{+12}$	0.76	1.62	$2.33 \cdot 10^{+12}$	0.79	3.3	1.30
20 (old)	$1.71 \cdot 10^{+12}$	0.36	2.4	$1.12 \cdot 10^{+12}$	0.38	5.6	1.47



**Figure 1.** Neutrino spectra for the  $\pi^+$  (left) and the  $\pi^-$  (right) focussed beam for a decay tunnel length of 60 m.

### 3. Sensitivity on $\theta_{13}$

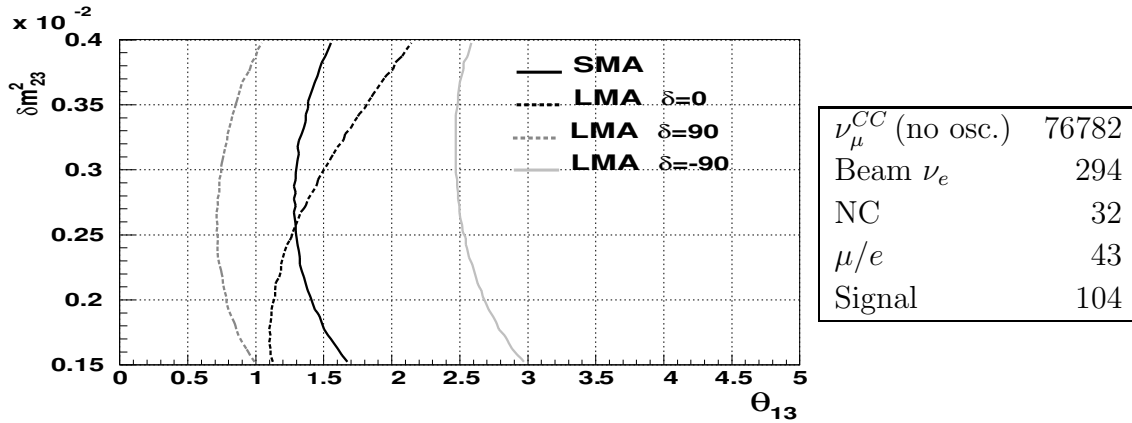
The  $\theta_{13}$  sensitivity is computed assuming  $\delta = 0$ , solar SMA solution,  $\delta m_{23}^2 = 2.5 \cdot 10^{-3}$  eV<sup>2</sup>,  $\theta_{23} = 45^\circ$  and 5 years of data taking. These are the standard benchmark assumptions used by similar projects [6], [10].

Fig. 2 shows the  $\theta_{13}$  sensitivity (90% CL) in case of no signal and summarizes the event rate computed for  $\theta_{13} = 2^\circ$ . The experiment would have sensitivity down to  $\theta_{13} = 1.2^\circ$  ( $\sin^2 2\theta_{13} = 1.75 \cdot 10^{-3}$ )

### 4. CP sensitivity

CP sensitivity is computed assuming a 2 year run with the  $\pi^+$  focussed beam and 8 years with the  $\pi^-$  focussed beam. This sharing is motivated by the unfavorable cross section ratio  $\bar{\nu}_e / \nu_e \sim 1/6$  at 300 MeV.

A 10% error on the solar  $\delta m^2$  and  $\sin^2 2\theta$ , as expected from the KamLAND experiment [11] and a 2% error on the atmospheric  $\delta m^2$  and  $\sin^2 2\theta$ , as expected from the JHF neutrino experiment [6] are taken into account. Correlations between  $\theta_{13}$  and



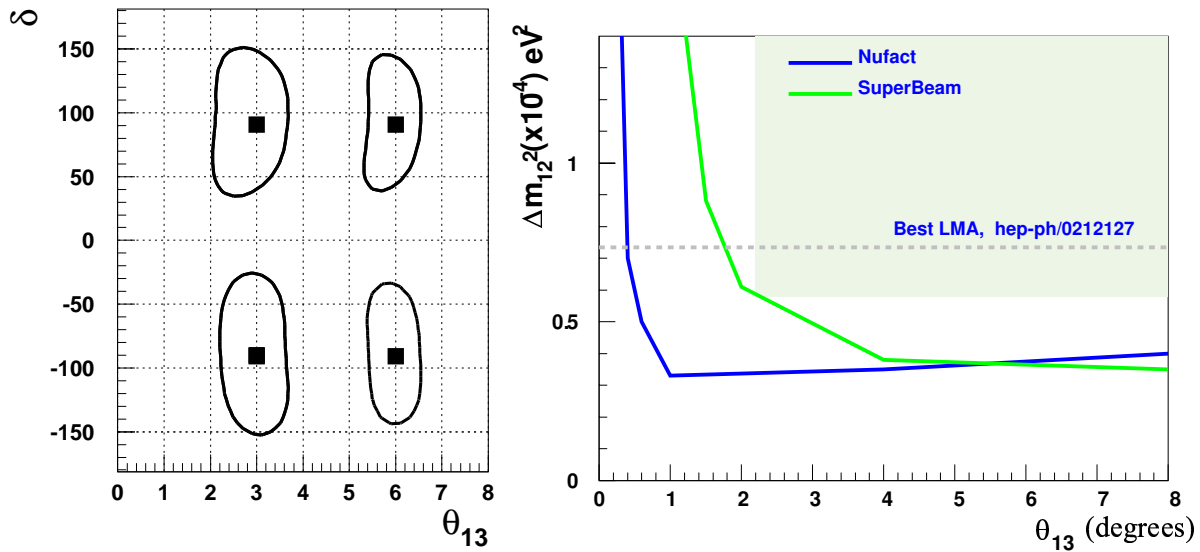
**Figure 2.** Left:  $\theta_{13}$  sensitivity (90%CL) computed for a 2200 kton exposure, under the solar SMA solution or under LMA and different  $\delta$  values. Right: Number of events for the same exposure, SMA solution, in case of  $\theta_{13} = 2^\circ$ .

$\delta$  are fully accounted for, while the sign  $\delta m_{13}^2$  and the  $\theta_{23}/(\pi/2 - \theta_{23})$  ambiguities are not considered. A systematic error of 2% is accounted for the signal efficiency and background normalization, as discussed in [2].

Solutions for different values of  $\delta$  and  $\theta_{13}$ , Fig.3-left, show very small correlations between the two parameters.

Since the sensitivity to CP violation heavily depends on the true value of  $\delta m_{12}^2$  and  $\theta_{13}$ , we prefer to express the CP sensitivity for a fixed value of  $\delta$  in the  $\delta m_{12}^2$ ,  $\theta_{13}$  parameter space. The CP sensitivity to separate  $\delta = 90^\circ$  from  $\delta = 0^\circ$  at the 99%CL as a function of  $\delta m_{12}^2$  and  $\theta_{13}$ , following the convention of [12], is plotted in Fig. 3-right.

It is fair to say that SPL SuperBeam CP sensitivity approaches the Neutrino Factory sensitivity in the parameter space that will be explored by the JHF experiment:  $\theta_{13} \geq 2.3^\circ$ .



**Figure 3.** Left:  $\theta_{13} - \delta$  fits (99% CL) computed for  $\delta m_{12}^2 = 10^{-4} \text{ eV}^2$ ,  $\sin^2 2\theta_{12} = 0.8$ . The squares indicate the starting points. Right: CP sensitivity of the SPL-SuperBeam, see text, compared with a 50 GeV Neutrino Factory producing  $2 \cdot 10^{20} \mu$  decays/straight section/year, and two 40 kton detectors at 3000 and 7000 km [12]; the shaded region corresponds to the allowed LMA solution and the  $\theta_{13}$  sensitivity of JHF.

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